

CONSIDERATIONES
DE THEORIA MOTVS LUNAE PERFICIEN-
DA ET IMPRIMIS DE EIVS
VARIATIONE.

Auctore

L. E V L E R O.

I.

Et si Theoria motuum Lunae a praestantissimis Geometris summo studio est inuestigata, atque adeo a Celeb. Professore Göttingensi *Mayero* Tabulae Lunares observationibus apprime satisficientes sunt in medium allatae, plurimum tamen adhuc abest, quo minus ipsa Theoria penitus excolta existimari possit. Quanquam enim forma istarum Tabularum ex Theoria est deriuata, quae etiam plures inaequalitates in motu Lunae accurate suppeditauit nonnullae tamen maximi momenti occurrunt, quarum quantitas ex solis observationibus est definita cum earum determinatio per solam Theoriam nimis incerta relinqueretur. Quin etiam nullum est dubium quin verus Lunae motus multo pluribus inaequalitatibus, quam quae in his Tabulis assignantur, perturbetur quae etsi in usu practico ob paruitatem facile praetermitti possunt, tamen in Theoria minime contemnendae videntur neque Theoria ante satis excolta

culta censei poterit, quam omnes prorsus motus inaequalitates, ne minimis quidem exceptis, accurate assignare valuerimus.

II. Ad Theoriam autem motuum Lunae feliciter inuestigandam, non statim ab eius motu vero exordium videtur, quemadmodum ab iis, qui hoc opus susceperunt, est factum; verus enim motus, quatenus non solum secundum longitudinem, sed etiam secundum latitudinem continuo perturbatur, tot tantisque difficultatibus implicatur, et penitus obruitur, vt singulis expediendis neque nostrae neque Analyseos vires sufficiant. Quam ob causam in hoc tam difficili negotio methodum ab Astronomis praecipue felicissimo cum successu vsitatam adhiberi conueniet, vt ante quam veros Lunae motus inuestigemus, casus nobis fingamus simpliciore, multo paucioribus difficultatibus obnoxios, quos si expedire licuerit, tum demum studia nostra continuo propius ad veritatem applicare licebit.

III. Primo igitur motus Lunae in latitudinem prorsus remouendus videtur, ita vt non huius, sed alius cuiusdam Lunae, quae in ipso eclipticae plano moueatur, motus sit inuestigandus; quandoquidem hoc modo calculus a grauissimis illis difficultatibus, quibus motus nodorum et inclinatio ad eclipticam premitur, liberatur. Deinde ne ipse solis motus quatenus non est vniformis molestiam facessat, hoc quoque obstaculum in principio tollatur,

et motus solis quasi esset vniformis spectetur. Hac ratione aliae inaequalitates inuestigandae non supererunt, nisi quae partim ab excentricitate orbitae lunaris, partim ab elongatione Lunae a Sole pendent. Ac si simplicitas adhuc maior desideretur, etiam excentricitas abiiciatur, et eiusmodi Lunae motus indagetur, quae sine vlla excentricitate in plano eclipticae moueretur sole cursum suum vniformiter absolvente. Hunc tantum casum adeo simplicem qui accurate et ad computum accommodate euoluere potuerit, is certe iam plurimum in Theoria praestitisse esset censendus.

Tab. I.

Fig. 2.

IV. Remota ergo inclinatione orbitae Lunaris, centrum terrae vt quiescens spectetur in T, et tabula referente planum eclipticae, sit tempore quodam t elapso centrum Lunae in L et Solis in S. Assumta iam recta fixa TA ad principium scilicet arietis ducta vocentur distantiae: $TL=v$, $TS=u$ et $LS=z$, et anguli $ATL=\Phi$, $ATS=\theta$, sitque breuitatis gratia $STL=\Phi-\theta=\eta$, erit $x=\sqrt{(uu-2vu\cos.\eta+vv)}$ vbi quidem distantia v est valde parua prae u . Porro demisso ab L in rectam TA perpendicularo LV sit $TV=x$ et $VL=y$, eritque $x=v\cos.\Phi$ et $y=v\sin.\Phi$. Hinc $x\cos.\Phi+y\sin.\Phi=v$ et $x\sin.\Phi-y\cos.\Phi=0$: Ergo differentiando $dx\cos.\Phi+dy\sin.\Phi-d\Phi(x\sin.\Phi-y\cos.\Phi)=dv$ seu $dx\cos.\Phi+dy\sin.\Phi=dv$ et $dx\sin.\Phi-dy\cos.\Phi+d\Phi(x\cos.\Phi+y\sin.\Phi)=0$ seu $dx\sin.\Phi-dy\cos.\Phi=-v d\Phi$. Porro denuo differentiando:

 ddx

$$\begin{aligned}
 ddx \operatorname{cof}.\Phi + ddy \operatorname{fin}.\Phi - d\Phi(dx \operatorname{fin}.\Phi - dy \operatorname{cof}.\Phi) &= ddv \text{ seu} \\
 ddx \operatorname{cof}.\Phi + ddy \operatorname{fin}.\Phi &= ddv - vd\Phi \\
 ddx \operatorname{fin}.\Phi - ddy \operatorname{cof}.\Phi + d\Phi(dx \operatorname{cof}.\Phi + dy \operatorname{fin}.\Phi) &= -dvd\Phi \\
 &= -vdd\Phi \text{ seu} \\
 ddy \operatorname{cof}.\Phi - ddx \operatorname{fin}.\Phi &= 2dvd\Phi + vdd\Phi.
 \end{aligned}$$

V. Iam massae Solis, terrae ac Lunae designentur litteris S, T et L, ita ut sint vires acceleratrices, quibus Luna vrgetur ad terram secundum $LT = \frac{T}{v}$, et ad solem secundum $LS = \frac{S}{z}$, quae ducta recta SL ipsi TS parallela resolvitur in has binas vires:

1°. Secundum $LT = \frac{Sv}{z^2}$ et 2°. secundum $LS = \frac{Su}{z^2}$.

Quia deinde terra ad solem vrgetur vi secundum $TS = \frac{S}{u}$, et ad Lunam vi secundum $TL = \frac{L}{v}$, hae vires contrarie in Lunam translatae dant vim secundum $Lt = \frac{S}{uu}$, et secundum $LT = \frac{L}{vv}$ ita ut iam Luna his viribus vrgeri censenda sit;

1°. Sec. $LT = \frac{T+L}{vv} + \frac{Sv}{z^2}$; 2°. sec. $Lt = \frac{S}{uu} - \frac{Su}{z^2}$

quae porro secundum directiones coordinatarum TV et VL, seu ducta LR ipsi TA parallela secundum LR et VL, resolutae dant

secundum LR vim $= \frac{T+L}{vv} \operatorname{cof}.\Phi + \frac{Sv}{z^2} \operatorname{cof}.\Phi + \frac{S}{uu} \operatorname{cof}.\theta - \frac{Su}{z^2} \operatorname{cof}.\theta$

secundum LV vim $= \frac{T+L}{vv} \operatorname{fin}.\Phi + \frac{Sv}{z^2} \operatorname{fin}.\Phi + \frac{S}{uu} \operatorname{fin}.\theta - \frac{Su}{z^2} \operatorname{fin}.\theta.$

Q 2

VI.

VI. His viribus inuentis fumendo temporis elemento dt constante, principia motus praebent has aequationes

$$\frac{d^2x}{dt^2} = -\frac{(T+L)\cos\Phi}{vv} - \frac{Sv\cos\Phi}{z^3} - \frac{S\cos\theta}{uu} + \frac{Su\cos\theta}{z^3}$$

$$\frac{d^2y}{dt^2} = -\frac{(T+L)\sin\Phi}{vv} - \frac{Sv\sin\Phi}{z^3} - \frac{S\sin\theta}{uu} + \frac{Su\sin\theta}{z^3}$$

vnde ob $ddy\cos\Phi - ddx\sin\Phi = 2dvd\Phi + vdd\Phi$

$$\text{et } ddx\cos\Phi + ddy\sin\Phi = ddv - vd\Phi^2$$

nanciscimur has binas aequationes principales:

$$1^\circ. \frac{2dvd\Phi + vdd\Phi}{dt^2} = \frac{S\sin\eta}{uu} - \frac{Su\sin\eta}{z^3}$$

$$2^\circ. \frac{ddv - vd\Phi^2}{dt^2} = -\frac{(T+L)}{vv} - \frac{Sv}{z^3} - \frac{S\cos\eta}{uu} + \frac{Su\cos\eta}{z^3}$$

Vt iam pro dt^2 valorem determinatum introducamus, consideremus motum Solis, qui cum ad terram follicitari censendus sit vi $\frac{S+T}{uu}$, habebitur simili modo:

$$\frac{2dud\theta + udd\theta}{dt^2} = 0 \text{ et } \frac{ddu - ud\theta^2}{dt^2} = -\frac{(S+T)}{uu}$$

sumamus iam Solis distantiam a terra mediam $= a$, et motum medium temporis t conuenientem $= \zeta$, erit ex posteriori aequatione $\frac{a d\zeta^2}{dt^2} = \frac{S+T}{aa}$, vnde colligimus $\frac{1}{dt^2} = \frac{T+S}{a^3 d\zeta^2}$

ficque loco elementi dt introducimus elementum cognitum pariter constans $d\zeta$, et has formulas adpiscimur:

$$1^\circ. 2dvd\Phi + vdd\Phi = \frac{S a^3 d\zeta^2 \sin\eta}{S+T} \left(\frac{1}{uu} - \frac{u}{z^3} \right)$$

$$2^\circ. ddv - vd\Phi^2 = -\frac{(T+L)a^3 d\zeta^2}{(T+S)vv} - \frac{S a^3 d\zeta^2}{S+T} \left(\frac{v}{z^3} + \frac{\cos\eta}{uu} - \frac{u\cos\eta}{z^3} \right)$$

vbi

vbi notandum est loco $\frac{S}{S+T}$ unitatem scribi licere cum massa T prae S evanescat.

VII. Vt litteras maiusculas S, T, L ex calculo exterminemus, contemplemur etiam motum Lunae medium, qui quidem esset futurus, si vires perturbantes a Sole oriundae abessent; hoc casu statuatur distantia Lunae media a terra $=c$, et ratio eius motus medii ad motum medium Solis $=n:1$; cum igitur sit $v=c$ et $d\Phi = nd\zeta$, posterior aequatio praebet $cnd\zeta^2 = \frac{(T+L)q^2 d\zeta^2}{(T+S)cc}$ vnde fit $\frac{T+L}{T+S} = \frac{nc^2}{a^2}$; ex quo nostrae aequationes principales sequentes induent formas:

$$1^\circ. 2dv d\Phi + v dd\Phi = a^2 d\zeta^2 \sin.\eta \left(\frac{1}{u} - \frac{u}{z^2} \right)$$

$$2^\circ. ddv - vd\Phi^2 = -\frac{nm^2}{vv} d\zeta^2 - \frac{v^2}{z^3} d\zeta^2 - a^2 d\zeta^2 \cos.\eta \left(\frac{1}{u} - \frac{u}{z^2} \right).$$

Totum ergo negotium huc redit, vt istae aequationes commode tractentur, ac si fieri queat ad integrationem perducantur: vbi quidem notasse iuuabit, membra posteriora quantitates u et z inuoluentia prae reliquis esse valde parua, indeque rationem approximandi esse petendam.

VIII. Ponamus autem breuitatis gratia:

$$\frac{1}{uu} - \frac{u}{z^2} = -M \text{ et } \frac{v}{z^3} + \cos.\eta \left(\frac{1}{u} - \frac{u}{z^2} \right) = N$$

vt aequationes nostrae fiant

$$1^\circ. 2dv d\Phi + v dd\Phi = -a^2 M d\zeta^2 \sin.\eta \text{ et}$$

$$2^\circ. ddv - vd\Phi^2 = -\frac{nm^2}{vv} d\zeta^2 - a^2 N d\zeta^2$$

vbi ob v prae u valde paruum et $z = \sqrt{(uu - 2uv \cos \eta + vv)}$ erit per approximationem

$$\frac{1}{z^2} = \frac{1}{u^2} + \frac{3v}{u^3} \cos \eta - \frac{3vv}{2u^4} (1 - 5 \cos \eta^2) - \frac{5v^2}{2u^5} (3 \cos \eta - 7 \cos \eta^3) + \frac{15v^3}{8u^6} (1 - 14 \cos \eta^2 + 21 \cos \eta^4) \text{ etc.}$$

ideoque litterarum M et N valores prodibunt

$$M = \frac{3v}{u^3} \cos \eta - \frac{3vv}{2u^4} (1 - 5 \cos \eta^2) - \frac{5v^2}{2u^5} (3 \cos \eta - 7 \cos \eta^3) + \frac{15v^3}{8u^6} (1 - 14 \cos \eta^2 + 21 \cos \eta^4)$$

$$N = \frac{v}{u^2} (1 - 3 \cos \eta^2) + \frac{3vv}{2u^3} (3 \cos \eta - 5 \cos \eta^3) - \frac{v^2}{2u^4} (3 - 30 \cos \eta^2 + 35 \cos \eta^4) - \frac{5v^3}{8u^5} (15 \cos \eta - 70 \cos \eta^3 + 63 \cos \eta^5)$$

vbi singula membra sequentia prae antecedentibus sunt vehementer exigua.

IX. Prima aequationum nostrarum ad integrabilitatem perducitur multiplicando eam per v tum vero etiam per $2v^2 d\Phi$, posteriori modo prodit $v^4 d\Phi^2 = -2a^2 d\zeta^2 fM v^3 d\Phi \sin \eta$.

Deinde prior multiplicetur per $2vd\Phi$ et posterior per $2dv$ ac summa dabit;

$$2dvddv + 2vdvd\Phi^2 + 2vvd\Phi dd\Phi = -2a^2 Mvd\zeta^2 d\Phi \sin \eta - \frac{2nac^2 dv}{v} d\zeta^2 - 2a^2 Nd\zeta^2 dv$$

vnde per integrationem eruitur:

$$dv^2 + vv d\Phi^2 = \frac{2nac^2 d\zeta^2}{v} - 2a^2 d\zeta^2 fMvd\Phi \sin \eta - 2a^2 d\zeta^2 fNdv$$

Statuamus breuitatis gratia:

$$a^2 fM v^3 d\Phi \sin \eta = -c^4 P \quad \text{et} \quad a^2 fM v d\Phi \sin \eta + a^2 fN dv = -ccQ$$

vt

vt obtineamus has formas:

$$1^{\circ}. v^4 d\Phi^2 = +2c^4 Pd\zeta^2 \text{ et } 2^{\circ}. dv^2 + vvd\Phi^2 = \frac{2nncc^3 d\zeta^2}{v} + 2ccQd\zeta^2$$

quae facta $v = cx$ fiunt:

$$1^{\circ}. x^4 d\Phi^2 = +2Pd\zeta^2 \text{ et } 2^{\circ}. dx^2 + xxxd\Phi^2 = \frac{2nn d\zeta^2}{x} + 2Qd\zeta^2$$

eritque:

$$dP = -\frac{a^2}{c} Mx^3 d\Phi \sin.\eta \text{ et } dQ = -\frac{a^2}{c} (Mxd\Phi \sin.\eta + Ndx)$$

existente

$$M = \frac{3ccx}{u^3} \cos.\eta - \frac{3ccxx}{2u^4} (1 - 5 \cos.\eta^2) - \frac{5c^3x^3}{2u^5} (3 \cos.\eta - 7 \cos.\eta^3)$$

$$N = \frac{c^3x}{u^3} (1 - 3 \cos.\eta^2) + \frac{3ccxx}{2u^4} (3 \cos.\eta - 5 \cos.\eta^3) - \frac{c^3x^3}{2u^5} (3 - 30 \cos.\eta^2 + 35 \cos.\eta^4).$$

X. Ex priore aequatione iam est $d\Phi = \frac{d\zeta \sqrt{2P}}{xx}$, qui valor in altera substitutus dat:

$$dx^2 + \frac{2Pd\zeta^2}{xx} = \frac{2nn d\zeta^2}{x} + 2Qd\zeta^2$$

vnde elicitur:

$$dx = d\zeta \sqrt{(2Q + \frac{2nn}{x} - \frac{2P}{xx})} \text{ vel etiam}$$

$$\frac{dx \sqrt{2P}}{xx} = d\Phi \sqrt{(2Q + \frac{2nn}{x} - \frac{2P}{xx})}$$

hincque discimus quantitatem $2Q + \frac{2nn}{x} - \frac{2P}{xx}$ nunquam fieri posse negativam; euanescere autem potest, quod fit dum Luna vel in apogeo versatur vel in perigeo, quandoquidem utroque casu fit $dx = 0$. Ceterum si vires perturbantes abessent pro motu medio,

medio, quo $x=1$ et $d\Phi=nd\zeta$ foret $n=V_2P$, seu $P=\frac{1}{2}nn$ et $nn=2nn+\frac{1}{2}Q$ seu $Q=-\frac{1}{2}nn$, qui ergo valores his litteris proxime conueniunt.

XI. Nisi excentricitas orbitae euanescat vel sit quam minima, eius introductio in calculum satis commode ad formulas differentiales primi gradus manuducit, quae ad computum astronomicum maxime videntur accommodatae. Duplici imprimis modo haec reductio institui potest, vnde deinceps alias latius patentes eiusmodi resolutiones deriuare licet. Alterum quidem modum iam alibi fusius sum persecutus, sed dignitas materiae omnino requirere videtur, vt vtrumque hic dilucide exponam, simulque cognationem ostendam, quo facilius intelligi possit, quanta emolumenta inde expectare liceat.

Reductio prior formularum inuentarum ope excentricitatis facta.

XII. Ordiamur a formula posteriori, quae per V_2 diuisa est:

$$\frac{d^2x}{x^2} V_2 P = d\Phi V_2 \left(Q + \frac{nn}{x} - \frac{P}{xx} \right)$$

ac statuamus $\frac{1}{x} = \frac{1-q \cos \omega}{p}$, seu $x = \frac{p}{1-q \cos \omega}$, vbi sequentia sunt obseruanda:

1°. Quantitas p in e ducta ob $v=cx$ exprimit semiparametrum orbitae, quatenus ea cum ellipfi comparatur; foretque p quantitas constans, si vires per-

perturbantes abessent, nunc autem erit quantitas variabilis.

2°. Quantitas q in eadem comparatione denotat excentricitatem, quae ob vires perturbantes pariter vt variabilis est spectanda.

3°. Angulus ω designat anomaliam veram ab apogeo computatam, et ob $v = cx$, erit distantia apogei $= \frac{cp}{1-q}$ et distantia perigei $= \frac{cp}{1+q}$, vnde semiaxis transuersus orbitae $= \frac{cp}{1-q^2}$.

4°. Loco vnus variabilis x introduximus tres nouas p , q et ω , inter quas autem iam vnam determinationem stabiliuimus qua dx euanescere debet si $\sin. \omega = 0$; alteram determinationem consideratio formulae irrationalis suppeditabit.

XIII. In formula $Q + \frac{nn}{x} - \frac{P}{xx}$ loco $\frac{x}{x}$ substituamus valorem assumptum $\frac{1-q \cos. \omega}{p}$, et prodibit

$$Q + \frac{nn}{p} - \frac{P}{pp} - \frac{nnq}{p} \cos. \omega + \frac{2Pq}{pp} \cos. \omega - \frac{Pqq}{pp} \cos. \omega^2$$

cuius radix quadrata quia factorem habere debet $\sin. \omega$

oportet vt sit 1°. $P = \frac{1}{2}nnp$, et 2°. $Q + \frac{nn}{p} - \frac{P}{pp} = \frac{Pqq}{pp}$

ficque prodeat $Q + \frac{nn}{x} - \frac{P}{xx} = \frac{Pqq}{pp} \sin. \omega^2$. Quo facto

erit $\frac{dx}{xx} \sqrt{P} = -\frac{qd\Phi \sin. \omega}{p} \sqrt{P}$, seu $\frac{dx}{xx} = -\frac{qd\Phi}{p} \sin. \omega$.

Cum autem sit $\frac{dx}{xx} = \frac{dp}{pp} + \cos. \omega \frac{dq}{p} - \frac{q}{p} d\omega \sin. \omega$,

habebimus $\frac{q}{p} (d\Phi - d\omega) \sin. \omega = -\frac{dp}{pp} - \cos. \omega \frac{dq}{p}$.

Ex factis autem binis hypothefibus erit primo

$p = \frac{2P}{nn}$ et ob $nn = \frac{2P}{p}$ altera dat $Q + \frac{P(1-q^2)}{pp} = 0$,

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feu $Q + \frac{n^2}{2p}(1 - qq) = 0$ hincque $\frac{1 - qq}{p} = -\frac{2Q}{n^2}$. Denique prima aequatio $d\Phi = \frac{d\zeta \sqrt{2P}}{xx}$ dat $\frac{d\Phi}{d\zeta} = \frac{n(1 - q \cos \omega)^2}{p \sqrt{p}}$, feu $d\zeta = \frac{p d\Phi \sqrt{p}}{n(1 - q \cos \omega)^2}$.

XIV. Quia nunc P et Q sunt quantitates, quarum differentialia saltem vt cognita spectantur, variationes momentaneae elementorum motus sequenti modo se habebunt.

1°. Pro quantitate p erit $dp = \frac{2 dP}{n^2}$, ideoque

$$dp = \frac{2 a^2}{n^2 c} M x^2 d\Phi \sin \eta \text{ vbi } x = \frac{p}{1 - q \cos \omega}$$

2°. Pro semiaxe orbitae $\frac{c p}{1 - qq}$ habemus $d \frac{1 - qq}{p} = -\frac{2 dQ}{n^2}$ ideoque $d \frac{1 - qq}{p} = \frac{2 a^2}{n^2 c} (M x d\Phi \sin \eta + N dx)$

quia vero est $dx = \frac{-q x d \cos \omega}{p} \sin \omega$ erit

$$d \frac{1 - qq}{p} = \frac{2 a^2}{n^2 c} x d\Phi (M \sin \eta - \frac{N q \sin \omega}{1 - q \cos \omega})$$

3°. Inuento differentiali quantitatis $\frac{1 - qq}{p}$, quam tantisper vocabo R, erit $qq = 1 - pR$ et $\frac{qq}{p^2} = \frac{1}{p} - \frac{R}{p}$, vnde fit

$$d \frac{qq}{p^2} = \frac{2 q}{p} d \frac{q}{p} = -\frac{2 dp}{p^3} + \frac{R dp}{p^2} - \frac{1}{p} dR = -\frac{(1 + qq) dp}{p^3} - \frac{1}{p} dR$$

vbi si loco dp et dR valores inuenti substituantur, reperitur

$$\frac{2 q}{p} d \frac{q}{p} = \frac{2 a^2 q x d\Phi}{n^2 c p} \left(\frac{M(2 \cos \omega + q \sin \omega^2) \sin \eta}{(1 - q \cos \omega)^2} + \frac{N \sin \omega}{1 - q \cos \omega} \right) \text{ ideoque}$$

$$d \frac{q}{p} = \frac{a^2 x d\Phi}{n^2 c} \left(\frac{M(2 \cos \omega + q \sin \omega^2) \sin \eta}{(1 - q \cos \omega)^2} + \frac{N \sin \omega}{1 - q \cos \omega} \right)$$

vnde

vnde concluditur :

$$\frac{q}{p}(d\Phi - d\omega)\sin.\omega = \frac{a^2}{n} \frac{a^2}{n c} x d\Phi \cdot \frac{M \sin.\eta}{(1 - q \cos.\omega)^2} - \cos.\omega d \cdot \frac{q}{p} \text{ feu}$$

$$\frac{q}{p}(d\Phi - d\omega)\sin.\omega = \frac{a^2 x d\Phi}{n n c} \left(\frac{M(2 \sin.\omega^2 - q \sin.\omega^2 \cos.\omega) \sin.\eta}{(1 - q \cos.\omega)^2} - \frac{N \sin.\omega \cos.\omega}{1 - q \cos.\omega} \right)$$

sicque habebimus :

$$d\Phi - d\omega = \frac{a^2 x x d\Phi}{n n c q} \left(\frac{M(2 - q \cos.\omega) \sin.\eta \sin.\omega}{1 - q \cos.\omega} - N \cos.\omega \right)$$

vnde motus lineae absidum definitur.

4^o. His variationibus definitis erit tandem

$$x = \frac{p}{1 - q \cos.\omega} \text{ et } d\zeta = \frac{p d\Phi \vee p}{n(1 - q \cos.\omega)^2}$$

qua posteriori formula ratio inter $d\Phi$ et $d\zeta$, illinc vero ratio inter $d\Phi$ et $d\omega$ exprimitur.

Reductio altera formularum inuentarum ad differentialia primi gradus.

XV. Aequationi posteriori haec inducatur forma :

$$\frac{d x}{x} \vee P = -d\Phi \vee (Q x x + n n x - P)$$

priore existente $x x d\Phi = d\zeta \vee 2 P$, et excentricitas ita introducatur vt ponatur $x = p + q \cos.\omega$, ficque distantia maxima fit $= p + q$ et minima $= p - q$, vbi autem quantitates p et q sunt variables. Cum nunc sit :

$\frac{d x}{x} = \frac{d p + d q \cos.\omega - q d \omega \sin.\omega}{p + q \cos.\omega}$, quae expressio evanescere debet si $\sin.\omega = 0$, valor ipsius x in altera parte substitutus dabit

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Q x x

$$Qxx + nnx - P = Qpp + 2Qpq \cos \omega + Qqq \cos \omega^2 \\ + nnp + nnq \cos \omega \\ - P.$$

Hic ergo ponatur $2Qp + nn = 0$ et $Qqq = -Qpp - nnp + P$ vt fiat $V(Qxx + nnx - P) = \sin \omega V(Qpp + nnp - P) = +q \sin \omega V - Q$. At ob $nnp = -2Qpp$, habemus $Qqq = Qpp + P$, feu

$$Q = \frac{-P}{pp - qq} = \frac{-nn}{2p}, \text{ vnde fit } \frac{nn(pp - qq)}{p} = 2P, \text{ et } \frac{nn}{p} = -2Q.$$

Quare altera aequatio hanc induit formam:

$$\frac{d\omega}{\omega} V P = -q d\Phi \sin \omega V - Q, \text{ feu } \frac{d\omega}{\omega} V(pp - qq) = -q d\Phi \sin \omega$$

vnde colligimus:

$$dp + dq \cos \omega - q d\omega \sin \omega = \frac{-q(p + q \cos \omega) d\Phi \sin \omega}{\sqrt{(pp - qq)}}.$$

Hinc singularum quantitatum variationes momentaneas ex differentialibus cognitis dP et dQ assignare poterimus.

1°. Aequatio $\frac{nn}{p} = -2Q$ dat $\frac{nn dp}{pp} = 2dQ$, ideoque

$$dp = \frac{2pp dQ}{nn}.$$

2°. Ex aequatione $\frac{nn(pp - qq)}{p} = 2P$ feu $p - \frac{qq}{p} = \frac{2P}{nn}$, sequitur

$$dp + \frac{qq dp}{pp} - \frac{2q dq}{p} = \frac{2dP}{nn}, \text{ feu } q dq = \frac{dp(pp + qq)}{2p} - \frac{pdP}{nn}$$

vnde fit $q dq = \frac{p(pp + qq)dQ - pdP}{nn}$.

3°. Hi valores in vltima aequatione substituti dabunt:

$$\frac{2pp dQ}{nn} + \frac{p(pp + qq)dQ \cos \omega}{nnq} - \frac{pdP \cos \omega}{nnq} - q d\omega \sin \omega = \frac{-q(p + q \cos \omega) d\Phi \sin \omega}{\sqrt{(pp - qq)}}$$

vnde

vnde fit:

$$qqd\omega \sin.\omega = \frac{p dQ}{nn} (2pq + (pp + qq) \cos.\omega) - \frac{p dP}{nn} \cos.\omega \\ + \frac{qq(p + q \cos.\omega) d\Phi \sin.\omega}{\sqrt{(pp - qq)}}$$

4°. Cum autem fit $dx = -\frac{qxd\Phi \sin.\omega}{\sqrt{(pp - qq)}}$ erit $dP = \frac{-a^3}{c} Mx^3 d\Phi \sin.\eta$
 et $dQ = \frac{-a^3}{c} d\Phi (Mx \sin.\eta - \frac{Nq x \sin.\omega}{\sqrt{(pp - qq)}})$, qui valores in
 formulis inuentis substituti praebeant:

$$dp = \frac{-2a^3 ppx d\Phi}{nnc} (M \sin.\eta - \frac{Nq \sin.\omega}{\sqrt{(pp - qq)}})$$

$$dq = \frac{a^3 pxd\Phi}{nnc} (M \sin.\eta (2p \cos.\omega - q \sin.\omega^2) + \frac{N(pp + qq) \sin.\omega}{\sqrt{(pp - qq)}})$$

$$d\omega = \frac{x d\Phi}{\sqrt{(pp - qq)}} - \frac{a^3 px d\Phi}{nncq} (M \sin.\eta (2p + q \cos.\omega) \sin.\omega \\ - \frac{N(2pq + (pp + qq) \cos.\omega)}{\sqrt{(pp - qq)}}).$$

Denique ob $2P = \frac{nn(pp - qq)}{p}$ est $dZ = \frac{xx d\Phi \sqrt{p}}{n \sqrt{(pp - qq)}}$ existente
 $x = p + q \cos.\omega$.

Reductio generalior binas praecedentes in se
 complectens.

XVI. Statuamus $x = \frac{p + q \cos.\omega}{1 - r \cos.\omega}$, vbi angulus ω
 ita se habeat vt casu $\sin.\omega = 0$ euanescat dx ; seu vt
 distantia fiat maxima casu $\omega = 0$, minima vero casu
 $\omega = 180^\circ$.

Erit ergo $\frac{dx}{x} = \frac{dp + dq \cos.\omega - q d\omega \sin.\omega}{p + q \cos.\omega} + \frac{dr \cos.\omega - r d\omega \sin.\omega}{1 - r \cos.\omega}$ seu

$$\frac{dx}{x} = \frac{dp + dq \cos.\omega}{p + q \cos.\omega} + \frac{dr \cos.\omega}{1 - r \cos.\omega} - \frac{(p + q) d\omega \sin.\omega}{(p + q \cos.\omega)(1 - r \cos.\omega)}$$

Nunc fiat substitutio in expressione $Qxx + nnx - P$
 quae abibit in hanc formam:

R 3

+ Qpp

$$\left. \begin{array}{l} +Qpp + 2Qpq \cos. \omega + Qqq \cos. \omega^2 \\ +nnP + nnq \cos. \omega - nnqr \cos. \omega^2 \\ -P - nnpr \cos. \omega - Pr r \cos. \omega^2 \\ + 2Pr \cos. \omega \end{array} \right\} (1 - r \cos. \omega)^2$$

Hic primo flatuatur $nn(pr - q) = 2Pr + 2Qpq$, deinde fit $Qpp + nnP - P = -Qqq + nnqr + Pr r$; vt fiat

$$\sqrt{(Qxx + nnx - P)} = \frac{\sqrt{(Qpp + nnP - P)}}{1 - r \cos. \omega} \sin. \omega \text{ ideoque}$$

$$\frac{dx}{x} \sqrt{P} = \frac{-d\Phi \sqrt{(Qpp + nnP - P)}}{1 - r \cos. \omega} \sin. \omega,$$

vnde sequentes determinaciones deducuntur.

XVII. Quaeramus primo rationem inter P et Q quae ob $nn = \frac{2Pr + 2Qpq}{pr - q}$ ex aequatione

$$Qpp + nnP - P = -Qqq + nnqr + Pr r$$

ita reperitur:

$$Q(p^2r - pqqr + ppq - q^2) = P(-pr + pr^2 - q + qrr)$$

quae per $pr + q$ diuisa dat

$$Q(pp - qq) = -P(1 - rr) \text{ seu } Q = \frac{-P(1 - rr)}{pp - qq}.$$

Hinc prior determinatio $nn(pr - q) = 2Pr + 2Qpq$ praebet

$$nn(pr - q) = 2Pr - \frac{2Ppq(1 - rr)}{pp - qq} = \frac{2P(ppr - qqr - pq + qrr)}{pp - qq}$$

et per $pr - q$ diuidendo $nn = \frac{2P(p + qr)}{pp - qq}$, ita vt fit

$$\frac{pp - qq}{p + qr} = \frac{2P}{nn} \text{ et } \frac{1 - rr}{p + qr} = \frac{-2Q}{nn}.$$

Deinde

Deinde loco nn iterum scribendo $\frac{2Pr+2Qpq}{pr-q}$ fit

$$Qpp + nnp - P = \frac{(pr+q)(P+Qpq)}{pr-q} = \frac{P(pr+q)^2}{pp-qq}$$

vnde concludimus :

$$\frac{dx}{x} = \frac{-(pr+q)d\Phi \sin.\omega}{(1-r \cos.\omega)\sqrt{(pp-qq)}}$$

At ob $x = \frac{p+q \cos.\omega}{1-r \cos.\omega}$, forma differentialis $\frac{dx}{x}$ ita exhiberi potest vt fit

$$\frac{dx}{x} = \frac{dp+dq \cos.\omega}{x(1-r \cos.\omega)} + \frac{dr \cos.\omega}{1-r \cos.\omega} - \frac{(pr+q)d\omega \sin.\omega}{x(1-r \cos.\omega)^2} = \frac{-(pr+q)d\Phi \sin.\omega}{(1-r \cos.\omega)\sqrt{(pp-qq)}}$$

Quocirca erit

$$\frac{(pr+q)d\omega \sin.\omega}{1-r \cos.\omega} = \frac{(pr+q)x d\Phi \sin.\omega}{\sqrt{(pp-qq)}} + dp + dq \cos.\omega + x dr \cos.\omega.$$

XVIII. Quodsi iam formulas superiores ad P et Q reductas differentiemus, ad sequentes expressiones perueniemus :

$$+ dp(pp+2pqr+qq) - dq(2pq+ppr+qqr) - qdr(pp-qq) \\ = \frac{2(p+qr)^2 dP}{nn.}$$

$$dp(1-rr) + rdq(1-rr) + qdr(1+rr) + 2prdr = \frac{2(p+qr)^2 dQ}{nn.}$$

vnde cum differentialia dP et dQ dentur, bina tantum trium elementorum dp , dq et dr definiuntur, tertio quasi arbitrio nostro relicto. Verum

$$\text{ob } dx = \frac{-(pr+q)x d\Phi \sin.\omega}{(1-r \cos.\omega)\sqrt{(pp-qq)}}$$

$$\text{erit : } dP = \frac{-a^2 x d\Phi}{c}. M x x \sin.\eta \text{ et}$$

$$dQ = \frac{-a^2 x d\Phi}{c} \left(M \sin.\eta - \frac{N(pr+q) \sin.\omega}{(1-r \cos.\omega)\sqrt{(pp-qq)}} \right).$$

Vel

Vel etiam angulum ω pro lubitu assumere licet, ac tum binis illis aequationibus hanc tertiam iungendo

$$dp + dq \cos \omega + x dr \cos \omega = \frac{(pr+q)d\omega \sin \omega}{1-r \cos \omega} - \frac{(pr+q)x d\Phi \sin \omega}{\sqrt{pp-qq}}$$

omnia tria elementa dp , dq et dr definiri poterunt. Denique ob $2P = \frac{nn(pp-qq)}{p+qr}$ erit $dZ = \frac{xx d\Phi \sqrt{p+qr}}{n \sqrt{pp-qq}}$.

XIX. Mirum videbitur, quod in hac reductione angulus ω arbitrio nostro relinquatur, cum certe positio et motus lineae absidum minime a nostra voluntate pendeant. Verum hic perpendi oportet, eatenus tantum distantiam $v = cx$ fieri maximam vel minimam factis $\sin \omega = 0$, quatenus idem angulus ω non in reliquis quantitates ita ingreditur, ut in valore pro $\frac{dx}{x}$ inuento factor $\sin \omega$ iterum tollatur. Quodsi exempli causa reperiretur $\sqrt{pp-qq} = s \sin \omega$, minime amplius concludere liceret posito $\sin \omega = 0$, formulam $\frac{dx}{x d\Phi}$ esse evanituram. Quocirca angulus ω nequaquam inter quantitates assumtas admitti potest, nisi forte constet a cuiusmodi angulo positio lineae absidum pendeat.

XX. Antequam hunc casum deseram, binas illas aequationes differentiales pro elementis dp , dq et dr inventas diligentius examinasse iuuabit. Ac si inde primo elementum dp elidatur reperitur:

$$\frac{dq(1-rr)(pr+q)}{p+qr} + dr(pr+q) = \frac{-(1-rr)dP + (pp+qq+2pqr)dQ}{nn}$$

sin autem inde elementum dq exterminetur, prodit

$$\frac{dp(1-rr)(pr+q)}{p+qr} + \frac{dr(pr+q)^2}{p+qr} = \frac{r(1-rr)dP + (2pq+ppr+qqr)dQ}{nn}$$

Eiecto

Eiecto autem elemento dr obtinetur

$$dp(pr+q) - \frac{dq(pr+q)^2}{p+qr} = \frac{(2pr+q+qrr)dP + (pp-qq)dQ}{nn}$$

Quod si iam harum binas quasque in locum illarum substituamus, calculus haud parum fiet simplicior hae vero videntur commodissimae:

$$dp - \frac{dq(pr+q)}{p+qr} = \frac{(2pr+q+qrr)dP + q(pp-qq)dQ}{nn(pr+q)}$$

$$dr + \frac{dq(1-rr)}{p+qr} = \frac{(1-rr)dP + (pp+qq+2pqr)dQ}{nn(pr+q)}$$

Vnde assumpto q reliqua elementa facile determinantur sin autem angulus ω vt cognitus spectetur, hinc valores pro dp et dr in postrema aequatione differentiali supra data (XVIII.) substituti determinationem elementi dq suppeditabunt. Peruenitur autem ad hanc aequationem:

$$\frac{dq(pr+q)\sin.\omega^2}{p+qr} = (pr+q)\sin.\omega \left(d\omega - \frac{d\Phi(p+q\cos.\omega)}{\sqrt{(pp-qq)}} \right)$$

$$- \frac{dP}{nn(pr+q)} (2pr+q+qrr - (p+qr)(1+rr)\cos.\omega - q(1-rr)\cos.\omega^2)$$

$$- \frac{dQ}{nn(pr+q)} (q(pp-qq) + (p+qr)(pp+qq)\cos.\omega + q(pp+qq+2pqr)\cos.\omega^2).$$

XXI. Substituendo denique hic pro dP et dQ valores supra indicatos (XVIII.)

$$\frac{dq\sin.\omega}{p+qr} = d\omega - \frac{d\Phi(p+q\cos.\omega)}{\sqrt{(pp-qq)}}$$

$$+ \frac{a^3 M \times d\Phi\sin.\eta\sin.\omega}{nnc(pr+q)(1-r\cos.\omega)^2} (2pp-qq+pqrr + (p+qr)(3q-pr)\cos.\omega - q(pr-q+2qrr)\cos.\omega^2)$$

$$- \frac{a^3 N \times d\Phi}{nnc(pr+q)(1-r\cos.\omega)} \left(\frac{q(pp-qq) + (p+qr)(pp+qq)\cos.\omega + q(pp+qq+2pqr)\cos.\omega^2}{\sqrt{(pp-qq)}} \right)$$

Tom. XIII. Nou. Comm.

S

si qui-

si quidem nunc totam aequationem per $(pr+q)\sin\omega$ diuidere licuit; commode enim vsu venit, vt membrum elemento $M\sin\eta$ affectum factorem $1-\cos\omega^2 = \sin\omega^2$ fortiretur.

Quodsi iam hic ponatur $q=0$, reductio resultat prior scribendo q loco r , sin autem ponatur $r=0$, reductio habetur posterior, vnde intelligitur quanto latius pateat haec reductio generalior ambas praecedentes in se complectens. Loco dP et dQ etiam in praecedentibus formulis substituantur valores ac reperietur:

$$\begin{aligned} dp &= \frac{dq(pr+q)}{p+qr} - \frac{a^3 M x d\Phi \sin\eta}{nnc(1-r\cos\omega)^2} (2pp-qq+pq r+2q(p+qr)\cos\omega \\ &\quad + q(pr+q)\cos\omega^2) \\ &\quad + \frac{a^3 N x d\Phi \sin\omega}{nnc(1-r\cos\omega)} q \sqrt{pp-qq} \\ dr &= \frac{-dq(1-rr)}{p+qr} - \frac{a^3 M x d\Phi \sin\eta}{nnc(1-r\cos\omega)^2} (pr+q-2(p+qr)\cos\omega+(pr-q \\ &\quad + 2qrr)\cos\omega^2) \\ &\quad + \frac{a^3 N x d\Phi \sin\omega}{nnc(1-r\cos\omega)} \left(\frac{pp+qq+2pqr}{\sqrt{pp-qq}} \right) \end{aligned}$$

et loco dq valorem superiorem substituendo:

$$\begin{aligned} \frac{dp \sin\omega}{p+qr} &= \frac{pr+q}{p+qr} \left(d\omega - \frac{d\Phi(p+q\cos\omega)}{\sqrt{pp-qq}} \right) \\ &\quad - \frac{a^3 M x d\Phi \sin\eta \sin\omega \cos\omega}{nnc(1-r\cos\omega)^2} (pr-q+2qrr\cos\omega) - \frac{a^3 N x d\Phi \cos\omega}{nnc(1-r\cos\omega)} \\ &\quad \left(\frac{pp+qq+2pqr}{\sqrt{pp-qq}} \right) \\ \frac{dr(pr+q)\sin\omega}{p+qr} &= - \frac{(pr+q)(1-rr)}{p+qr} \left(d\omega - \frac{d\Phi(p+q\cos\omega)}{\sqrt{pp-qq}} \right) \\ &\quad - \frac{a^3 M x d\Phi \sin\eta \sin\omega}{nnc(1-r\cos\omega)^2} (2p+qr-prr+(q-3pr-3qrr \\ &\quad + pr^2)\cos\omega+r(pr-q+2qrr)\cos\omega^2) \\ &\quad + \end{aligned}$$

$$+ \frac{a^3 N x d\Phi}{nnc(1-r\cos\omega)\sqrt{(pp-qq)}} (ppr + 2pq + qqr + (1-rr)(pp+qq)\cos\omega - r(pp+qq+2pqr)\cos\omega^2).$$

XXII. Si excentricitas orbitae satis fuerit notabilis, commodissime reductione prima utemur, quia ibi aberrationes a motu regulari in ellipfi facto definiuntur. Sin autem excentricitas fuerit quam minima vel adeo nulla, neque primam reductionem neque secundam in usum vocare licebit, quandoquidem anomaliae ω tum ne locus quidem relinquatur; ac spectata quantitate q saltem ut minima, quia ea denominatorem formulae pro $d\Phi - d\omega$ inventae afficit, motus lineae absidum nimis fit vagus et incertus. Neque etiam adhuc perspicio, quomodo postrema reductio sumendo $\omega = \eta$ in hac investigatione utilitatem afferre posset, tam propter multitudinem, quam complicationem formularum, quas resolvere oporteret. Nihilo tamen minus casus quo excentricitas plane evanesceret sine dubio pro simplicissimo esset habendus; ex quo in eius resolutione merito omne studium collocandum videtur quo his difficultatibus superatis deinceps veri motus lunaris investigatio feliciori successu suscipi, neque tantum ad usum practicum satis convenienter, sed etiam multo accuratius absolvi queat. Neque autem ad hunc casum evoluendum alia via aptior videtur, quam ut ad ipsas aequationes differentio-differentiales revertamur indeque approximationes idoneas petamus.

Inuestigatio motus si Luna in ecliptica sine vlla excentricitate sol autem vniformiter moueretur.

XXIII. Ponamus in ipsis aequationibus differentio differentialibus $v = cx$, et habebimus.

$$1^{\circ}. 2 dx d\Phi + x dd\Phi + \frac{a^3}{c} M d\zeta^2 \text{ fin. } \eta = 0$$

$$2^{\circ}. ddx - x d\Phi^2 + \frac{n n}{x x} d\zeta^2 + \frac{a^3}{c} N d\zeta^2 = 0$$

et quia motus solis assumitur vniformis erit $u = a$ et $\theta = \zeta$ ideoque $\Phi = \zeta + \eta$, hinc

$$\frac{a^3}{c} M = 3x \text{ cof. } \eta - \frac{3c x x}{2a} (1 - 5 \text{ cof. } \eta^2) - \frac{5c c x^3}{2a a} (3 \text{ cof. } \eta - 7 \text{ cof. } \eta^3)$$

$$\frac{a^3}{c} N = x(1 - 3 \text{ cof. } \eta^2) + \frac{3c c x x}{2a} (3 \text{ cof. } \eta - 5 \text{ cof. } \eta^3) - \frac{c c x^3}{2a a} (3 - 30 \text{ cof. } \eta^2 + 35 \text{ cof. } \eta^4)$$

vnde binae nostrae aequationes erunt

$$1^{\circ}. \left\{ \begin{array}{l} \frac{2 dx d\Phi}{d\zeta^2} + \frac{x dd\Phi}{d\zeta^2} \\ + 3x \text{ fin. } \eta \text{ cof. } \eta - \frac{3c}{2a} x x \text{ fin. } \eta (1 - 5 \text{ cof. } \eta^2) - \frac{5c c}{2a a} x^3 \text{ fin. } \eta (3 \text{ cof. } \eta - 7 \text{ cof. } \eta^3) \end{array} \right\} = 0$$

$$2^{\circ}. \left\{ \begin{array}{l} \frac{ddx}{d\zeta^2} - \frac{x d\Phi^2}{d\zeta^2} + \frac{n n}{x x} \\ + x(1 - 3 \text{ cof. } \eta^2) + \frac{3c c x x}{2a} (3 \text{ cof. } \eta - 5 \text{ cof. } \eta^3) - \frac{c c}{2a a} x^3 (3 - 30 \text{ cof. } \eta^2 + 35 \text{ cof. } \eta^4) \end{array} \right\} = 0$$

vbi cum $\frac{c}{a}$ fit quantitas quam minima, has aequationes in partes sectas concipere licet, quae sequentibus multo sint maiores, ad quem ordinem etiam approximationem accommodari conuenit.

XXIV. Si omnis perturbatio abesset, foret ob excentricitatem evanescentem, vti vidimus, $x=1$ et $\frac{d\phi}{d\xi} = n$ hincque $\frac{d\eta}{d\xi} = n-1$. Nunc perturbatione accedente statuamus:

$$x = 1 + P + Q + R \text{ et } \frac{d\phi}{d\xi} = n + p + q + r$$

hincque $\frac{d\eta}{d\xi} = n-1 + p + q + r$, vbi P, Q, R et p, q, r series maxime decrescentes referant, cum seriis superioribus ex perturbatione natis comparandas ac has ipsas quantitates tanquam functiones anguli η spectemus, siquidem nouimus omnes inaequalitates ab hoc solo angulo pendere. Erit ergo $dx = dP + dQ + dR$ et per $\frac{d\eta}{d\xi} = (n-1) + p + q + r$ multiplicando:

$$\frac{d^2x}{d\xi^2} = \left\{ \begin{array}{l} (n-1)dP + (n-1)dQ + (n-1)dR \\ \quad + p dP \quad + \quad p dQ \\ \quad \quad \quad + q dP \end{array} \right\} : d\eta$$

quae forma differentiatâ sumto iam elemento $d\eta$ constante dabit

$$\frac{d^2x}{d\xi^2} = \left\{ \begin{array}{l} (n-1)ddP + (n-1)ddQ + (n-1)ddR \\ \quad + p ddP \quad + \quad p ddQ \\ \quad + dP dp \quad + \quad dp dQ \\ \quad \quad \quad + q ddP \\ \quad \quad \quad + dq dP \end{array} \right\} : d\eta$$

multiplicetur denuo per $\frac{d\eta}{d\xi}$, prodibitque

$$\frac{d d \varpi}{d \xi^2} = \left\{ \begin{array}{l} (n-1)^2 ddP + (n-1)^2 ddQ + (n-1)^2 ddR \\ + 2(n-1)pddP + 2(n-1)pddQ \\ + (n-1)dpdP + (n-1)dpdQ \\ + 2(n-1)qddP \\ + (n-1)dq dP \\ + pp d d P \\ + pdpdP \end{array} \right\} : d\eta^2$$

simili modo cum fit

$\frac{d d \Phi}{d \xi^2} = dp + dq + dr$ per $\frac{d\eta}{d\xi}$ multiplicando erit

$$\frac{d d \Phi}{d \xi^2} = \left\{ \begin{array}{l} (n-1)dp + (n-1)dq + (n-1)dr \\ + pdp + pdq \\ + qdp \end{array} \right\} : d\eta$$

et $\frac{d \Phi^2}{d \xi^2} = nn + 2np + 2nq + 2nr$
 $+ pp + 2pq$ ac tandem

$$\frac{1}{x} = 1 - 2P - 2Q - 2R \\ + 3PP + 6PQ \\ - 4P^2.$$

XXV. Hos igitur valores in aequationes nostras introductos secundum ordines stabilitos distribuamus, vbi quidem elementum $d\eta$, quippe quod sponte intelligitur, omittamus.

* Aequa-

* Aequatio Prima

II.	III.	IV.
$+2n(n-1)dP$	$+2n(n-1)dQ$	$+2n(n-1)dR$
$+(n-1)dp$	$+2npdP$	$+2npdQ$
$+3\sin.\eta\cos.\eta$	$+2(n-1)p dP$	$+2nq dP$
	$+(n-1)dq$	$+2(n-1)p dQ$
	$+p dp$	$+2pp dP$
	$+(n-1)P dp$	$+2(n-1)q dP$
	$+3P\sin.\eta\cos.\eta$	$+(n-1)dr$
	$-\frac{3c}{2a}\sin.\eta(1-5\cos.\eta^2)$	$+p dq$
		$+q dp$
		$+(n-1)Pdq$
		$+Pp dp$
		$+(n-1)Qdp$
		$+3Q\sin.\eta\cos.\eta$
		$-\frac{3c}{a}P\sin.\eta(1-5\cos.\eta^2)$
		$-\frac{5cc}{2aa}\sin.\eta(3\cos.\eta-7\cos.\eta^3)$

Hic scilicet ordo primus deest, quia sublata perturbatione primae aequationis omnia membra sponte euanescunt.

Pro

* Huius aequationis nonnisi integrale particulare hic quaeritur, quod scilicet hypothesi assumtae, qua excentricitas euanescit, conveniat, et manifesto huiusmodi habet formam $P = A + B\cos.\eta^2$. Integrale autem completum foret

$$P = A + B\cos.\eta^2 + M\sin.\frac{n}{n-1}\eta + N\cos.\frac{n}{n-1}\eta$$

vbi M et N sunt constantes arbitrariae, quibus conditio excentricitatis continetur. Id quod peculiarem evolutionem meretur.

Pro sequentibus autem ordinibus terminos ad quemuis pertinentes seorsim nihilo aequari oportet.

XXVI. Aequatio altera sequenti modo in membra distribuitur.

Aequatio secunda.

I.	II.	III.	IV.
$-nn$	$(n-1)^2 ddP$	$+(n-1)^2 ddQ$	$+(n-1)^2 d d R$
$+nn$	$-2np$	$+2(n-1)pddP$	$+2(n-1)pddQ$
	$-3nnP$	$+(n-1)dpdP$	$+2(n-1)qddP$
	$+(1-3\text{cof.}\eta^2)$	$-2nq$	$+(n-1)dpdQ$
		$-pp$	$+(n-1)dqdP$
		$-2npp$	$+ppddP$
		$-3nnQ$	$+pdpdP$
		$+3nnPP$	$-2nr$
		$+P(1-3\text{cof.}\eta^2)$	$-2pq$
		$+\frac{c}{2a}(3\text{cof.}\eta-5\text{cof.}\eta^3)$	$-2nPg$
			$-Ppp$
			$-2nQq$
			$-3nnR$
			$+6nnPQ$
			$-4nnP^2$
			$+Q(1-3\text{cof.}\eta^2)$
			$+\frac{c}{a}P(3\text{cof.}\eta-5\text{cof.}\eta^3)$
			$-\frac{c^2}{2aa}(3-30\text{cof.}\eta^2+35\text{cof.}\eta^4)$

vbi membrum primum sponte se tollit.

XXVII. Secundus ordo ex vtraque aequatione quantitibus secundo loco assumtis definiendis infer-
vit,

vit, quae sunt P et p , ideoque ex his duabus aequationibus determinandae.

$$1^{\circ}. 2n(n-1)dP + (n-1)dp + 3d\eta \sin. \eta \cos. \eta = 0$$

$$2^{\circ}. (n-1)^2 ddP - 2npd\eta^2 - 3nnPd\eta^2 + d\eta^2(1-3\cos.\eta^2) = 0.$$

Prior autem integrata dat $2n(n-1)P + (n-1)p = \Delta + \frac{3}{2}\cos.\eta^2$
 feu $p = -2nP + \frac{\Delta}{n-1} + \frac{3\cos.\eta^2}{2(n-1)}$, qui valor in altera substitutus praebet:

$$(n-1)^2 ddP + nnPd\eta^2 - \frac{2n}{n-1}\Delta d\eta^2 - \frac{3n\eta^2 \cos.\eta^2}{n-1} + d\eta^2(1-3\cos.\eta^2) = 0$$

$$\text{feu } (n-1)^2 ddP + nnPd\eta^2 - \frac{2n}{n-1}\Delta d\eta^2 - \frac{3(2n-1)}{n-1}d\eta^2 \cos.\eta^2 + d\eta^2 = 0.$$

Statuamus, quandoquidem forma integralis sponte patet $P = A + B\cos.\eta^2$, si esset $nn = 4(n-1)^2$ poni deberet $P = A + B\cos.\eta^2 + C\eta \sin. \eta \cos. \eta$, erit $\frac{dP}{d\eta} = -2B\sin.\eta \cos. \eta$ et $\frac{d^2P}{d\eta^2} = -2B\cos.\eta^2 + 2B\sin.\eta^2 = 2B - 4B\cos.\eta^2$ et facta substitutione oritur:

$$\left. \begin{aligned} &+ 2(n-1)^2 B + nnA - \frac{2n}{n-1}\Delta + 1 \\ &- 4(n-1)^2 B\cos.\eta^2 + nnB\cos.\eta^2 - \frac{3(2n-1)}{n-1}\cos.\eta^2 \end{aligned} \right\} = 0$$

hincque $B = \frac{-3(2n-1)}{(n-1)(n-2)(3n-2)}$ et $\frac{2n}{n-1}\Delta = 1 + nnA + 2(n-1)^2 B$

et $p = -2nA - 2nB\cos.\eta^2$

$$+ \frac{1}{2n} + \frac{3}{2(n-1)} \cos.\eta^2.$$

$$+ \frac{1}{2}nA$$

$$+ \frac{(n-1)^2}{n} B.$$

Quare si ponamus:

$$P = A + B\cos.\eta^2 \quad \text{et} \quad p = \mathcal{A} + \mathcal{B}\cos.\eta^2$$

quantitas A arbitrio nostro relinquitur. eritque

$$B = \frac{-\frac{3}{2}(2n-1)}{(n-1)(n-2)(3n-2)} \text{ atque}$$

$$\mathfrak{A} = \frac{1}{2n} - \frac{3}{2}nA + \frac{(n-1)^2}{n}B \text{ et } \mathfrak{B} = \frac{3}{2(n-1)} - 2nB.$$

Quantitas A ideo manet indefinita, ut vel distantia media vel motus medius ad veritatem definiri possit, ob perturbationem enim, si c conueniat cum distantia media n non amplius cum ratione $\frac{d\Phi}{d\xi}$ congruit et vicissim.

XXVIII. Ad quantitates tertii ordinis Q et q determinandas, has habemus aequationes:

$$2n(n-1)dQ + (n-1)dq + 2(2n-1)p dP + p dp + (n-1)P dp + 3P \sin \eta \cos \eta - \frac{3c}{2a} \sin \eta (1-5 \cos \eta^2) = 0$$

$$(n-1)^2 ddQ - 2nq - 3nnQ + 2(n-1)p ddP + (n-1)dp dP - pp - 2nPp + 3nmPP + P(1-3 \cos \eta^2) + \frac{3c}{2a}(3 \cos \eta - 5 \cos \eta^3) = 0.$$

Cum autem sit $P = A + B \cos \eta^2$ et $p = \mathfrak{A} + \mathfrak{B} \cos \eta^2$ erit $dP = -2B \sin \eta \cos \eta$ et $dp = -2\mathfrak{B} \sin \eta \cos \eta$

hi valores in prima aequatione substituti dant

$$\frac{2n(n-1)dQ + (n-1)dq}{d\eta} - 4(2n-1)\mathfrak{A}B \sin \eta \cos \eta - 4(2n-1)\mathfrak{B}B \sin \eta \cos \eta^3 - \frac{3c}{2a} \sin \eta (1-5 \cos \eta^2) = 0$$

$$- 2 \mathfrak{A} \mathfrak{B}$$

$$- 2 \mathfrak{B} \mathfrak{B}$$

$$- 2(n-1)\mathfrak{B}A$$

$$- 2(n-1)\mathfrak{B}B$$

$$+ 3A$$

$$+ 3B$$

vnde

vnde per integrationem elicatur

$$2n(n-1)Q + (n-1)q + (2(2n-1)AB + 2B + (n-1)BA - \frac{1}{2}A) \text{ cof. } \eta^2 + \frac{3c}{2a} \text{ cof. } \eta - \frac{5c}{2a} \text{ cof. } \eta^3 = \Delta + ((2n-1)BB + \frac{1}{2}BB + \frac{1}{2}(n-1)BB - \frac{1}{4}B) \text{ cof. } \eta^4$$

fit breuitatis gratia

$$2(2n-1)AB + 2B + (n-1)BA - \frac{1}{2}A = \alpha$$

$$\frac{1}{2}(5n-3)BB + \frac{1}{2}BB - \frac{1}{4}B = \beta$$

$$\text{erit } q = -2nQ + \frac{\Delta}{n-1} - \frac{\alpha}{n-1} \text{ cof. } \eta^2 - \frac{\beta}{n-1} \text{ cof. } \eta^4 - \frac{3c}{2a(n-1)} \text{ cof. } \eta + \frac{5c}{2a(n-1)} \text{ cof. } \eta^3$$

Tum pro altera aequatione ob $ddP = 2B - 4B \text{ cof. } \eta^2$ est

$2(n-1)pddP$	$4(n-1)AB - 8(n-1)AB \text{ cof. } \eta^2 - 8(n-1)BB \text{ cof. } \eta^4$		
$+(n-1)dpdP$	$+4(n-1)BB$		
$-pp$	$-AA - 2AB$	$-4(n-1)BB$	
$-2nPP$	$-2nAA - 2nAB$	$-2nBB$	
$+3nnPP$	$-2nBA$		
$+P(1-3\text{cof. } \eta^2)$	$+3nnAA + 6nnAB + 3nnBB$		
	$+A - 3A - 3B$		
	$+B$		
$+\frac{(n-1)^2 ddQ}{a\eta^2}$	$-\frac{2n\Delta}{n-1} + \frac{2n\alpha}{n-1}$	$+\frac{2n\beta}{n-1}$	$+\frac{3c}{2a} \text{ cof. } \eta - \frac{5c}{2a} \text{ cof. } \eta^3$
$+nnQ$			$+\frac{37c}{a(n-1)} \text{ cof. } \eta - \frac{57c}{a(n-1)} \text{ cof. } \eta^3$

XXIX. Pro resolutione huius aequationis poni oportere manifestum est:

$$Q = C + D \text{ cof. } \eta^2 + E \text{ cof. } \eta^4 + F \text{ cof. } \eta + G \text{ cof. } \eta^3$$

$$\text{eritque } \frac{ddQ}{d\eta^2} = 2D - 4D \text{ cof. } \eta^2 - 16E \text{ cof. } \eta^4 - F \text{ cof. } \eta - 9G \text{ cof. } \eta^3$$

T 2

Hinc

Hinc istae nascuntur aequationes :

$$0 = n^2 C + 2(n-1)^2 D + 4(n-1) \mathfrak{A} B - 2 \mathfrak{A}^2 - 2n \mathfrak{A} A + 3m A A + A - \frac{2n \Delta}{n-1}$$

$$0 = 12(n-1)^2 E - (n-2)(3n-2) D - 2(5n-4) \mathfrak{A} B + 8(n-1) \mathfrak{B} B - 2 \mathfrak{A} \mathfrak{B} - 2n \mathfrak{B} A + 6m A B - 3A + B + \frac{2n \alpha}{n-1}$$

$$0 = -(3n-4)(5n-4) E - 12(n-1) \mathfrak{B} B - \mathfrak{B}^2 - 2n \mathfrak{B} B + 3m B B - 3B + \frac{2n \epsilon}{n-1}$$

$$0 = 6(n-1)^2 G + (2n-1) F + \frac{3(5n-3)c}{2(n-1)a}$$

$$0 = -(2n-3)(4n-3) G - \frac{5(5n-3)c}{2(n-1)a}$$

unde ordine retrogado litterae G, F, E et D determinantur tum vero ex prima valor ipsius Δ quaeratur, ut C maneat quantitas indefinita, ac tum etiam valor litterae q innotescet per quantitates iam definitas A, B, \mathfrak{A} et \mathfrak{B} , ex quibus α et ϵ resultant. Si enim ponatur $q = \mathfrak{C} + \mathfrak{D} \cos. \eta + \mathfrak{E} \cos. \eta^2 + \mathfrak{F} \cos. \eta + \mathfrak{G} \cos. \eta^2$ erit

$$\mathfrak{C} = -2n C + \frac{\Delta}{n-1}$$

$$\mathfrak{D} = -2n D - \frac{\alpha}{n-1}$$

$$\mathfrak{E} = -2n E - \frac{\epsilon}{n-1}$$

$$\mathfrak{F} = -2n F - \frac{3c}{2(n-1)a}$$

$$\mathfrak{G} = -2n G + \frac{5c}{2(n-1)a}$$

XXX. Simili modo perturbationes sequentium ordinum ex aequationibus supra datis colligi posse per

per se est manifestum. Calculus quidem haud parum sit molestus ac taediosus, verum sufficit methodum eum euoluendi hic dilucide exposuisse, ita ut nulla difficultas sit metuenda praeter calculi prolixitatem. Interim sequentia membra ita fiunt parva, ut pro usu astronomico facile reiici queant. Inuentis autem his omnibus litteris binae aequationes quibus iam motus lunae continetur ita se habent:

$$x = r + A + C + (B + D)\cos.\eta^2 + E\cos.\eta^4 + F\cos.\eta^6 + G\cos.\eta^8$$

$$\frac{d\phi}{d\zeta} = n + H + I + (J + K)\cos.\eta^2 + L\cos.\eta^4 + M\cos.\eta^6 + N\cos.\eta^8$$

existente $d\phi = d\eta + d\zeta$, et $\frac{d\eta}{d\zeta} = \frac{d\phi}{d\zeta} - 1$; vbi manifestum est constantem A nihilo aequalem poni posse, dummodo C in calculo retineatur, ut ratio media $d\phi : d\zeta$ ob sequentes terminos aliquantillum a vero valore n depulsa corrigi et ad veritatem reduci possit modo infra exponendo.

Operae autem pretium erit hos singulos terminos euoluere sumendo pro n valorem per observationes definitum, quia eadem inaequalitates; etiam si ad casum hunc maxime particularem pertinentes, tamen in vero quoque lunae motu locum inueniunt.

XXXI. Sumamus ergo $A = 0$, et cum motum lunae medium cum solis motu comparando sit $n = 13,25586$, calculus pro determinatione harum inaequalitatum ita se habebit:

T 3

 $n = 13,$

$n = 13,25586$	$l(2n-1) = 1,4067394$
$n-1 = 12,25586$	$l_3 = 0,4771213$
$2n-1 = 25,51172$	$l_3(2n-1) = 1,8838607$
$3n-2 = 37,76758$	$l(n-1) = 1,0883440$
$n-2 = 11,25586$	$l(n-2) = 1,0513787$
$B = -0,0146899$	$l(3n-2) = 1,5771191$
$\frac{x}{2(n-1)} = +0,1223904$	$3,7168418$
$2nB = -0,3894546$	$l-B = 8,1670189$
$\mathfrak{B} = +0,5118450$	$l^{\frac{3}{2}} = 0,1760913$
$\frac{x}{2n} = +0,0377191$	$l(n-1) = 1,0883440$
$\frac{(n-1)^2}{n} B = -0,1664559$	$l \frac{x}{2(n-1)} = 9,0877473$
$\mathfrak{A} = -0,1287368$	$l_2 = 0,3010300$
	$ln = 1,1224079$
	$l_2n = 1,4234379$
	$l-2nB = 9,5904568$
	$l(n-1)^2 = 2,1766880$
	$l \frac{(n-1)^2}{n} = 1,0542801$
	$l - \frac{(n-1)^2}{n} B = 9,2212990$

XXXII. Nunc pro litteris α et \mathfrak{E} calculus ita se habebit :

$2(n-1)\mathfrak{A}B = +0,096492$	$l\mathfrak{A} = -9,1097027$
$\mathfrak{A}\mathfrak{B} = -0,065893$	$lB = -8,1670189$
$\alpha = +0,030599$	$l_2 = 0,3010300$
$\frac{3}{2}(5n-3) = 31,63965$	$l(2n-1) = 1,4067394$
$\frac{3}{2}(5n-3)\mathfrak{B}\mathfrak{B} = -0,237897$	$+ 8,9844910$
$\frac{1}{2}\mathfrak{B}\mathfrak{B} = +0,130993$	$l\mathfrak{A} = -9,1097027$
$-0,106904$	$l\mathfrak{B} = +9,7091385$
	$- 8,8188412$
	$- \frac{3}{4}B$

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$$\begin{aligned}
 -\frac{1}{2}B &= +0,011018 & l\frac{1}{2}(5n-3) &= 1,5002316 \\
 E &= -0,095886 & lB &= +9,7091385 \\
 & & lB &= -8,1670189 \\
 & & & \underline{-9,3763890} \\
 & & lB B &= 9,4182770 \\
 & & l_2 &= 0,3010300 \\
 & & & \underline{9,1172470}
 \end{aligned}$$

Hinc primo quaeratur littera E.

$$\begin{aligned}
 -12(n-1)BB &= +1,105820 & lBB &= -7,8761574 \\
 -2nBB &= +0,199340 & l(n-1) &= 1,0883440 \\
 & + 1,305160 & l_{12} &= 1,0791812 \\
 -BB &= -0,261986 & & \underline{-0,0436826} \\
 & + 1,043174 & l_2 n &= 1,4234379 \\
 +3nnBB &= +0,113756 & & \underline{-9,2995953} \\
 -3B &= +0,044069 & lBB &= 6,3340378 \\
 & + 1,200999 & lnn &= 2,2448158 \\
 +\frac{2nE}{n-1} &= -0,207419 & l_3 &= 0,4771213 \\
 (3n-4)(5n-4)E &= +0,993580 & & \underline{+9,0559749} \\
 ergo E &= +0,00044603 & lE &= -8,9817552 \\
 & & l_2 n &= 1,4234379 \\
 & & & \underline{-0,4051931} \\
 & & l(n-1) &= 1,0883440 \\
 & & & \underline{-9,3168491} \\
 +0,99358 &= 9,9972028 \\
 l(3n-4) &= 1,5534895 \\
 l(5n-4) &= 1,7943437 \\
 & \underline{3,3478332} \\
 lE &= +6,6493696 \\
 & \text{XXXIII.}
 \end{aligned}$$

XXXIII. Porro pro littera D.

$12(n-1)^2 E = +0,803968$ $-2(5n-4)2B = -0,235557$ $+8(n-1)2B = -0,737223$ $-22B = +0,131786$ $B = -0,014690$ $+ \frac{3n\alpha}{n-1} = +0,066192$ $(n-2)(3n-2)D = +0,014476$ $\text{Ergo } D = +0,000034052$ $2(n-1)^2 D = +0,010230$ $+4(n-1)2B = +0,192984$ $-22 = -0,016573$ $nC - \frac{2n\Delta}{n-1} = +0,186641$ $\frac{\Delta}{n-1} = 0,007040 + \frac{1}{2}nC$ $\text{Ergo } C = 0,007040 - \frac{1}{2}nC$ $-2nD = -0,0009028$ $-\frac{\alpha}{n-1} = -0,0024970$ $\text{Ergo } \mathcal{D} = -0,0033998$ $-2nE = -0,0118250$ $-\frac{E}{n-1} = +0,0078237$	$1E = +6,6493696$ $l(n-1)^2 = 2,1766880$ $l12 = 1,0791812$ $+9,9052388$ $l22B = +7,5777516$ $l(5n-4) = 1,7943437$ $+9,3720953$ $l\alpha = +8,4857072$ $l\frac{2n}{n-1} = 0,3350939$ $+8,8208011$ $l0,014476 = 8,1606486$ $l(n-2)(3n-2) = 2,6284978$ $lD = 5,5321508$ $l(n-1)^2 = 2,1766880$ $l2 = 0,3010300$ $8,0098688$ $l22 = +8,2194054$ $l\dots = 9,2710070$ $l2n = 1,4234379$ $7,8475691$ $lD = 5,5321508$ $l2nD = +6,9555887$ $l\alpha = +8,4857072$ $l(n-1) = 1,0883440$ $+7,3973632$ $lE = +6,6493696$ $l2n = 1,4234379$ $+8,0728075$ <p style="text-align: right;">Ergo</p>
-----------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------	--------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------------

Ergo $\mathbb{C} = -0,0040013$

$l\mathbb{E} = -8,9817552$

$(2n-3)(4n-3)G = -12,90796 \frac{c}{a}$

$l(n-1) = 1,0883440$

$-7,8934112$

$l \frac{5n-3}{2} = 1,5002316$

$l(n-1) = 1,0883440$

$l \frac{5n-3}{2(n-1)} = 0,4118876$

$l5 = 0,6989700$

$1,1108576$

$G = -0,010975 \frac{c}{a}$

$1,1108576 = l12,90796$

$1,3712844 = l(2n-3)$

$1,6991735 = l(4n-3)$

$3,0704579$

$-8,0403997 = lG$

$2,1766880 = l(n-1)^2$

$0,7781513 = l6$

$-0,9952390$

$+6(n-1)^2G = -9,89097 \frac{c}{a}$

$\frac{2(5n-3)c}{2(n-1)a} = +7,74478 \frac{c}{a}$

$(2n-1)F = +2,14619 \frac{c}{a}$

$l \frac{5n-3}{2(n-1)} = 0,4118876$

$l3 = 0,4771213$

$0,8890089$

$l2,14619 = 0,3316683$

$l(2n-1) = 1,4067394$

$lF = +8,9249289$

$lG = -8,0403997$

$l2n = 1,4234379$

$l2nF = +0,3483668$

$l2nG = -9,4638376$

Ergo $F = +0,084126 \frac{c}{a}$

$-2nF = -2,23032 \frac{c}{a}$

$-\frac{3c}{2(n-1)a} = -0,122390 \frac{c}{a}$

$\mathbb{F} = -2,35271 \frac{c}{a}$

$-2nG = +0,29096 \frac{c}{a}$

$+\frac{5c}{2(n-1)a} = +0,20399 \frac{c}{a}$

$\mathbb{G} = +0,49495 \frac{c}{a}$

XXXIV. Ex his igitur valoribus colligimus:

$$x = 1 + C - 0,014656 \cos. \eta^2 + 0,000446 \cos. \eta^4 \\ + 0,084126 \frac{c}{a} \cos. \eta - 0,010975 \frac{c}{a} \cos. \eta^3 \\ \frac{d\Phi}{d\eta^2} = 13,134163 - \frac{2}{3}nC + 0,508445 \cos. \eta^2 - 0,004001 \cos. \eta^4 \\ - 2,3527 \frac{c}{a} \cos. \eta + 0,4949 \frac{c}{a} \cos. \eta^3$$

vbi constans C ita definiri debet, vt motus medius ex posteriori forma erutus praecise conueniat cum motu medio ex obseruationibus deducto. In hunc autem finem potestates $\cos. \eta^2$ et $\cos. \eta^4$ ad cosinus angulorum simplicium reduci debent, quia inde partes constantes emergunt cum principali coniungendae. Scilicet ob $\cos. \eta^2 = \frac{1}{2} + \frac{1}{2} \cos. 2\eta$ et $\cos. \eta^4 = \frac{3}{8} + \frac{1}{2} \cos. 2\eta + \frac{1}{8} \cos. 4\eta$, fit pars constans:

$$13,386885 - \frac{2}{3}nC \text{ ipsi } n = 13,25586 \text{ aequanda,} \\ \text{vnde fit } \frac{2}{3}nC = 0,131025, \text{ ideoque } C = 0,0065895.$$

Euolutis autem potestatibus $\cos. \eta$ reperitur

$$x = 0,999428 - 0,007105 \cos. 2\eta + 0,000056 \cos. 4\eta \\ + 0,07589 \frac{c}{a} \cos. \eta - 0,00274 \frac{c}{a} \cos. 3\eta \\ \frac{d\Phi}{d\eta^2} = 13,25586 + 0,252222 \cos. 2\eta - 0,000500 \cos. 4\eta \\ - 1,9815 \frac{c}{a} \cos. \eta + 0,1237 \frac{c}{a} \cos. 3\eta.$$

XXXV. Quo hinc facilius ipsum angulum Φ definire queamus, ponamus breuitatis gratia $\frac{d\Phi}{d\eta} = n+r$ erit

$$\frac{d\Phi}{d\eta} = \frac{n+r}{n-1+r} = \frac{n}{n-1} - \frac{r}{(n-1)^2} + \frac{rr}{(n-1)^3}$$

fit

fit $r = \alpha \cos. 2\eta + \xi \cos. 4\eta + \gamma \cos. \eta + \delta \cos. \eta^3$ erit

$rr = \frac{1}{2} \alpha \alpha + \frac{1}{2} \alpha \alpha \cos. 4\eta$ omiffis reliquis terminis, qui ad ordines fequentes deuoluerentur, et ob parvitatem facile negliguntur. Integratione ergo inflituta prodit

$$\Phi = \Delta + \frac{\pi}{n-1} \eta - \frac{\alpha \sin. 2\eta}{2(n-1)^2} - \frac{\xi \sin. 4\eta}{4(n-1)^2} - \frac{\gamma \sin. \eta}{(n-1)^2} - \frac{\delta \sin. 3\eta}{3(n-1)^2} \\ + \frac{\alpha \alpha}{2(n-1)^2} \eta + \frac{\alpha \alpha \sin. 4\eta}{8(n-1)^2}$$

vbi est:

$$\alpha = +0,252222; \xi = -0,000500; \gamma = -1,9815 \frac{c}{a}; \\ \delta = +0,1237 \frac{c}{a}$$

atque iam ante quidem C ita definiri debuiffet, vt et hic particula $\frac{\alpha \alpha}{2(n-1)^2} \eta$ tolleretur, prodiretque fecondum motum medium $\Phi = \Delta + \frac{\pi}{n-1} \eta = \Delta + 1,081593 \eta$. Singulis igitur terminis euolutis et in minuta fecunda conuerfis habebitur:

$$\Phi = \Delta + 1,081593 \eta - 173'', 177 \sin. 2\eta + 1'', 063 \sin. 4\eta \\ + 2721'' \frac{c}{a} \sin. \eta - 57'' \frac{c}{a} \sin. 3\eta.$$

XXXVI. Sed cum η ex motu medio non innotefcat, ratio primo inter ζ et η eft ftabilienda, quae ob $\Phi = \zeta + \eta$ elicitur:

$$\zeta = \Delta + \frac{\pi}{n-1} \eta - \frac{\alpha \sin. 2\eta}{2(n-1)^2} - \frac{\xi \sin. 4\eta}{4(n-1)^2} - \frac{\gamma \sin. \eta}{(n-1)^2} - \frac{\delta \sin. 3\eta}{3(n-1)^2} \\ + \frac{\alpha \alpha \sin. 4\eta}{8(n-1)^2}$$

V 2

hinc-

hincque colligitur :

$$\eta = \text{Const.} + 12,25586\zeta + 2122''\cdot 43 \sin. 2\eta - 13''\cdot 023 \sin. 4\eta \\ - 33348''\frac{c}{a} \sin. \eta + 694''\frac{c}{a} \sin. 3\eta$$

unde haud difficulter ad datam solis longitudinem mediam angulus η colligitur, tum vero erit $\Phi = \eta + \zeta$. Denique distantia lunae a terra habebitur :

$$v = c(0,999428 - 0,007105 \cos. 2\eta + 0,000056 \cos. 4\eta) \\ + 0,07589 \frac{c}{a} \cos. \eta - 0,00274 \frac{c}{a} \cos. 3\eta$$

At ex massis Solis, lunae, et terrae quantitas c ita definitur ut fit $\frac{n n c^3}{a^3} = \frac{T + L}{T + S} = \frac{T}{S}$, unde patet ob actionem solis distantiam lunae mediam aliquantulum imminui.

XXXVII. In vero lunae motu caedem istae inaequalitates quoque occurrunt, unde haud inutile erat eas omni cura determinasse; ab Astronomis autem nomine variationis lunae designantur, quia omnes in vna tabula comprehendi possunt, argumentum distantiae solis a luna prae se ferente. Patet autem eius partem posteriorem a parallaxi solis pendere, seu a fractione $\frac{c}{a}$, dum prior absolute datur. Quare si quantitas huius inaequalitatis pro variis angulis per observationes innotesceret, inde vicissim parallaxis solis concludi posset. Cum igitur Tabulae *Mayerianae* cum coelo ita exacte conveniant,

niant, vt inaequalitates tanquam ex obseruationibus conclusae spectari queant, comparatio nostrae formulae inuentae cum his Tabulis, parallaxin solis nobis exhibere poterit. Consideremus solum casum, quo angulus $\eta = 90^\circ$, quia tum pars variationis absoluta euanescit, eritque per formulam nostram variatio $= -34042'' \cdot \frac{a}{c}$ tabulae autem *Mayerianae* habent $-1', 57'' = -117''$ vnde sequitur $\frac{a}{c} = \frac{34042}{117} = 291$, cui rationi cum ratio parallaxium sit aequalis, parallaxis autem lunae media sit $57', 15'' = 3435''$, erit parallaxis solis $= \frac{3435}{291} = 11\frac{4}{5}''$. Haec fortasse methodus parallaxin solis definiendi reliquis excepto veneris transitu, longe anteferenda videtur, si quidem tabulae *Mayerianae* nunquam ultra minutum a coelo dissident, quia enim haec variationis portio ad $117''$ affurgit, leuis mutatio in parallaxi solis assumpta sensibilem aberrationem a veritate produceret vt scilicet parallaxis solis prodiret $= 8\frac{1}{2}''$ tabulae *Mayerianae* loco $-117''$ habere deberent $-84''$, ex hac autem solis parallaxi foret $\frac{a}{c} = 400$. Considerari potest quoque maxima variatio angulo $\eta = 135^\circ$ fere respondens, quae ex nostra forma est $= -2122'' - 23050'' \cdot \frac{a}{c}$; at ex Tabulis *Mayerianis* $= -41', 41'' = -2501''$, vnde sequitur $\frac{a}{c} = \frac{23050}{279}$, sed haec conclusio minus est certa, ob effectum a parallaxi solis ortum multo minorem. Contra vero maxima variatio hinc potius oriri videtur $= -2202''$ seu $36', 42''$. Verum hic probe animaduerti oportet,

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ex excentricitate partem quoque ipsi $\sin. 2\eta$ proportionalem nasci, quae in his tabulis cum vera variatione est coniuncta. Haecque est causa, cur parallaxin solis ex variatione ubi $\sin. 2\eta = 0$ et $\sin. 4\eta = 0$ feliciter determinare licuerit minime vero ex variatione maxima.

ANNO-